Dynamical Systems and Chaos Part II: Biology Applications

Lecture 10: Coupled Systems.

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Foreword

- ▶ In order to model "populations" of physical/biological systems, represented as dynamical systems, one has to introduce **coupling** between sub-systems.
- Such a population will be the dynamical system itself, i.e. numerically and mathematically sub-systems are not distinguished.
- ► Structure of the coupling is crucial (one-to-one, all-to-all coupling etc.).
- ▶ In the coupled systems new dynamical regimes arise.
- Novelty is mainly due to breaking of symmetry (in amplitudes and/or phase).
- ► Amplitude: homogeneity/non-homogeneity.
- ▶ Phase: synchrony/asynchrony.

One-to-one coupling

Neighboring elements get connected (e.g. neurons, harmonic oscillators, closely located pendulum clocks, neuro-mascular oscillators in the small intestine, coupling on a ring).



All elements of the system connect to each other (bacterial population, biofilms, cells in a tissue, embryo etc.).



Homogeneity/inhomogeneity

- Symmetry break in variables' steady state levels, oscillatory levels or amplitude.
- ▶ This applies both to steady states and oscillations.
- Generally, through the pitchfork bifurcation (breaking of symmetry).



Figure: Left: homogeneous oscillatory regime — two subsystems have the same amplitude and levels of oscillations. Right: two subsystems have different amplitude and oscillatory levels — inhomogeneous solution.

(A)synchrony

 First observation: Huygens, pendulum clocks hanging on a wooden beam (17th century).



Figure 10.1: Different types of synchronization.

E. Izhikevich, Dynamical Systems in Neuroscience, MIT Press, 2007.

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Phase of oscillation

- Let x_0 be an arbitrary point on a periodic orbit γ .
- Any other point on γ can be characterized by how much time θ has passed since the last visiting of x₀.
- ϑ is called *phase of oscillation*.
- The phase is always bounded by the period of oscillation T. Sometime ϑ is normalized by T.
- The phase is also defined outside of γ with *isochrons*, set of initial conditions having the same phase.
- Many oscillators can be rewritten in a simpler phase model form in vicinity of their γ sets:

$$\vartheta' = 1$$

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▶ It is useful to assume that the phase is defined on a unit circle (just as sine or cosine functions).

Phase of oscillation



E. Izhikevich, Dynamical Systems in Neuroscience, MIT Press, 2007.

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Isochrons



E. Izhikevich, Dynamical Systems in Neuroscience, MIT Press, 2007.

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- Phase and amplitude has different relative stability to the external perturbation (external forcing or other oscillators).
- If timescale of relaxation to the limit cycle is τ₁ and timescale of the coupling/interaction between oscillators is τ₂, weak coupling implies τ₁ ≪ τ₂.
- ► In this case, coupling does not affect amplitude of oscillation, but its phase ⇒ essence of the synchronization phenomena.

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Weak forcing: Arnold tongues

$$\frac{d\vartheta}{dt} = \omega_0 + \varepsilon Q(\vartheta, \omega t)$$

where ϑ and ω_0 are the oscillator's phase and frequency, respectively; Q is the coupling function, ω and ε are the external force frequency and amplitude, respectively.



Figure: Arnold tongues.

- ▶ Non-constant interaction.
- Single pulses at certain time moments (flashing fireflies, neurons).
- ► A dynamical system with a pulsed coupling can be represented:

$$\dot{x} = f(x) + A\delta(t - t_s)$$

where A is an amplitude of the signal given as a Dirac $\delta(t)$ function at time moment t_s .

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• One just resets a state variable x by the constant A.

Phase response curve



E. Izhikevich, Dynamical Systems in Neuroscience, MIT Press, 2007.

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Two coupled harmonic oscillators

$$\begin{cases} I \begin{cases} \frac{dx_1}{dt} = y_1 \\ \frac{dy_1}{dt} = -\frac{k}{m}x_1 + \frac{k_1}{m}(x_2 - x_1) \\ II \begin{cases} \frac{dy_2}{dt} = -\frac{k}{m}x_1 - \frac{k_1}{m}(x_2 - x_1) \end{cases} \end{cases}$$

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Two coupled harmonic oscillators: in- and anti-phase



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Two harmonic oscillators: more complex oscillations

• One can vary both initial conditions and frequencies.



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Two coupled Brusselators

$$\begin{cases} I \begin{cases} \frac{dx_1}{dt} = A - (B+1)x_1 + x_1^2 y_1 \\ \frac{dy_1}{dt} = Bx_1 - x_1^2 y_1 + D(y_2 - y_1) \\ \\ II \begin{cases} \frac{dx_2}{dt} = A - (B+1)x_2 + x_2^2 y_2 \\ \frac{dy_2}{dt} = Bx_2 - x_2^2 y_2 + D(y_1 - y_2) \end{cases} \end{cases}$$

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Homogeneous dynamical regimes



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In-homogeneous dynamical regimes



E. Volkov and V. Romanov, Physica Scripta, 51, 19–28, 1995.

Trying to find in-homogeneous steady state



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Trying to find out-of-phase oscillations



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In-homogeneous limit cycle



E. Volkov and V. Romanov, Physica Scripta, 51, 19–28, 1995.

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In-homogeneous limit cycle



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In-homogeneous limit cycle \rightarrow chaos



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- ▶ Global coupling, i.e. all-to-all.
- ▶ Similar to 2 coupled Brusselators.

$$\frac{dx_i}{dt} = A - (B+1)x_i + x_i^2 y_i
\frac{dy_i}{dt} = Bx_i - x_i^2 y_i + D(\frac{1}{N} \sum_{j=1}^N y_j - y_i)$$

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Wave-solution: out-of-phase homogeneous oscillations



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Other out-of-phase homogeneous oscillations



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Quorum Sensing (QS)



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${\rm Repressilator}\,+\,{\rm QS}$





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J. Garcia-Ojalvo, M. Elowitz, and S. Strogatz, *PNAS*, **101**:10955-10960, 2004.

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$$\begin{split} \frac{da_i}{dt} &= \frac{\alpha}{1+C_i^n} - a_i & \frac{dA_i}{dt} = \beta_a(a_i - A_i) \\ \frac{db_i}{dt} &= \frac{\alpha}{1+A_i^n} - b_i & \frac{dB_i}{dt} = \beta_b(b_i - B_i) \\ \frac{dc_i}{dt} &= \frac{\alpha}{1+B_i^n} - c_i + \kappa \frac{S_i}{1+S_i} & \frac{dC_i}{dt} = \beta_c(c_i - C_i) \\ & \frac{dS_i}{dt} = -k_{s0}S_i + k_{s1}A_i - \eta(S_i - Q\bar{S}) \,, \end{split}$$
where $\bar{S} = \frac{1}{N} \sum_{j=1}^N S_j$.

J. Garcia-Ojalvo, M. Elowitz, and S. Strogatz, PNAS, 101:10955-10960, 2004.

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I. Potapov, E. Volkov, and A. Kuznetsov, Phys Rev E, 031901, 2011.

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• β determines which comes first: HB_{in} or HB_{ant}.



I. Potapov, E. Volkov, and A. Kuznetsov, Phys Rev E, 031901, 2011.

- ▶ There are dynamical behaviors not predicted by the bifurcation analysis.
- ▶ In-homogeneous anti-phase solution changes its properties after the "torus" bifurcation (TR₂). The resulting solutions are not seen in the bifurcation analysis.



I. Potapov, E. Volkov, and A. Kuznetsov, Phys Rev E, 031901, 2011.

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Repressilator + QS: In-homogeneous limit cycle



I. Potapov, E. Volkov, and A. Kuznetsov, Phys Rev E, 031901, 2011.

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Repressilator + QS: after torus bifurcation (TR_2)



I. Potapov, E. Volkov, and A. Kuznetsov, Phys Rev E, 031901, 2011.

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Repressilator + QS: attractive vs. repulsive coupling



- The same equations except for the AI.
- <u>Phase-repulsive</u> coupling has AI equation:

$$\frac{dS_i}{dt} = -k_{s0}S_i + \underline{k_{s1}B_i} - \eta(S_i - Q\bar{S}),$$

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where
$$\bar{S} = \frac{1}{N} \sum_{j=1}^{N} S_j$$
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E. Ullner et al., Phys Rev E, 031904, 2008.

Phase-repulsive coupling: steady state diagram



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E. Ullner et al., Phys Rev E, 031904, 2008.

Phase-repulsive coupling: in-homogeneous oscillations



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E. Ullner et al., Phys Rev E, 031904, 2008.

Phase-repulsive coupling: anti-phase solution branch



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E. Ullner et al., Phys Rev E, 031904, 2008.

Phase-repulsive coupling: in-phase solution unstable



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E. Ullner et al., Phys Rev E, 031904, 2008.

Phase-repulsive coupling: quantitative analysis



E. Ullner et al., Phys Rev E, 031904, 2008.

Phase-repulsive coupling: chaos



E. Ullner et al., Phys Rev E, 031904, 2008.

Phase-repulsive coupling: variety of kinetics



E. Ullner et al., Phys Rev E, 031904, 2008.

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Summary

- Coupling between mathematically similar systems is plausible assumption of a population of "identical" (genetically, morphologically etc.) cells.
- ▶ Coupling increases the dynamical complexity of models, that is the dynamics of the coupled systems is much more complex (in many cases) than that of an isolated system.
- Coupling increases the mathematical complexity of models, since dimension grows with increasing number of coupled systems.
- ▶ Type of coupling is essential, for example, all-to-all (diffusion coupling, QS etc.), one-to-one (neurons), pulsed coupling...
- Dynamical behaviors in coupled systems are classified roughly by homogeneity and synchronization, that is amplitude and temporal characteristics are compared.

 Non-technical presentation of the synchronization phenomenon:

Coupled oscillators and Biological Synchronization by Steven Strogatz and Ian Stewart, in Scientific American, December 1993, p. 102.

 Scholarpedia pages on synchronization: http:

//www.scholarpedia.org/article/Synchronization

 Scholarpedia pages on phase models (more neuronal dynamics oriented):

http://www.scholarpedia.org/article/Phase_model